

Phenomenological thermodynamics of multiple conserved quantities

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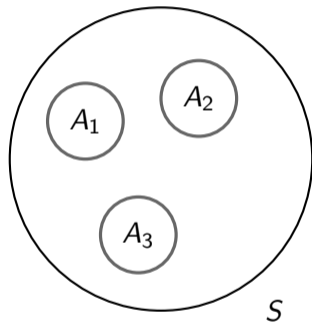
December 8, 2022

Why consider phenomenological thermodynamics

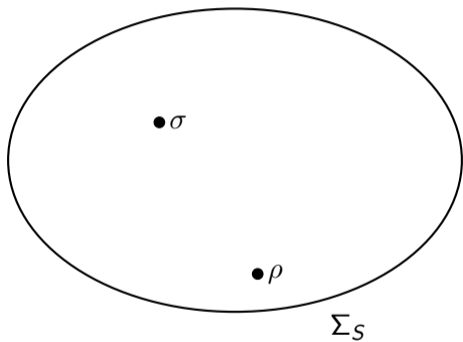
- ▶ Axiomatic framework makes assumptions clear¹.
- ▶ Multiple quantities make the role of information more explicit: Landauer's principle.
- ▶ No microscopic theory needed: Black holes.

¹Philipp Kammerlander. "Tangible Phenomenological Thermodynamics". en. PhD thesis. ETH Zurich, 2019. DOI: 10.3929/ETHZ-B-000413414, Philipp Kammerlander and Renato Renner. "Tangible phenomenological thermodynamics". In: (Feb. 20, 2020). arXiv: 2002.08968 [math-ph]. URL: <https://arxiv.org/abs/2002.08968>.

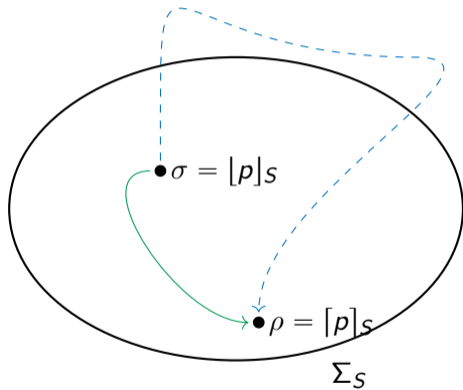
A **system** is a non-empty and finite subset of the thermodynamic world \mathcal{S}



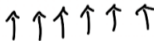
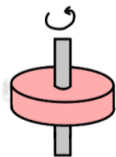
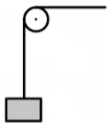
The **state space** of the system is Σ_S



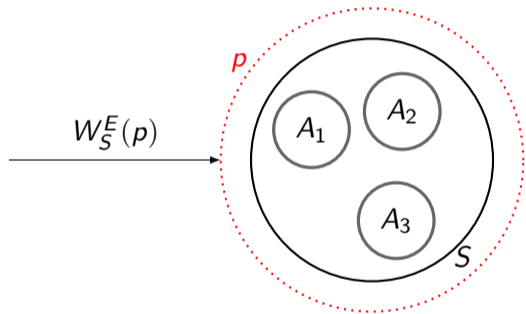
Process $p \in \mathcal{P}$



Work = Quantity + Information

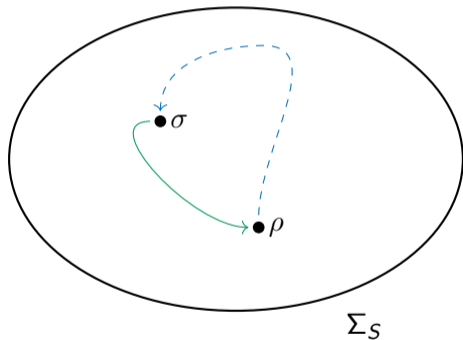


$W_S^B(p)$: work of quantity B done by process p on system S .



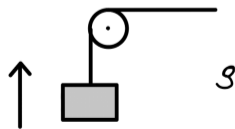
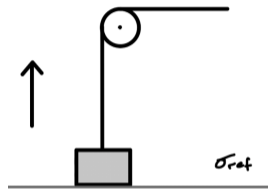
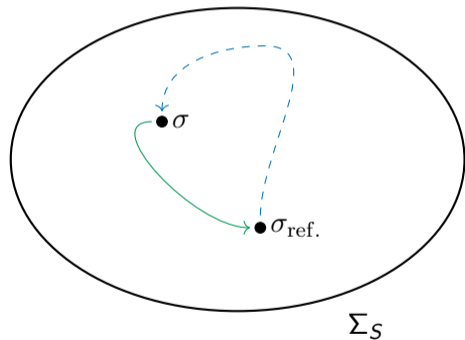
First law

- ▶ There is a non-dissipative process connecting any two states in Σ_S .



Internal B -Worth

- ▶ For p a work process on a system S $U_S^B(\sigma) = +/ - W_S^B(p) + U_S^B(\sigma_{\text{ref.}})$



Heat = quantity

$$Q_S^B(p) = \Delta U_S^B - W_S^B(p)$$

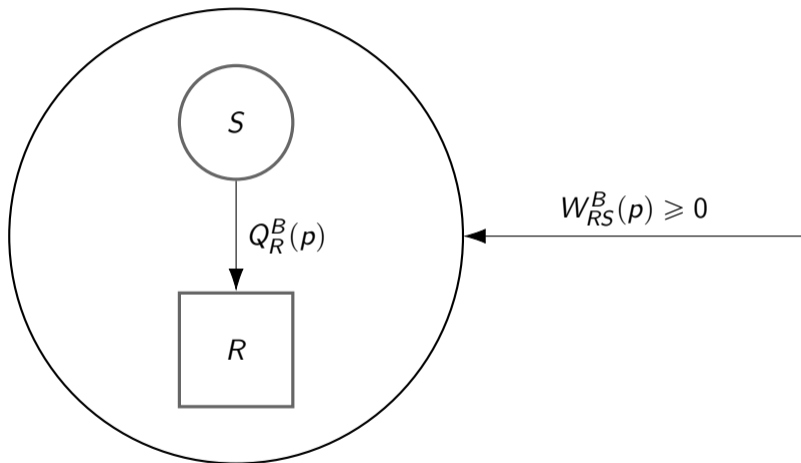


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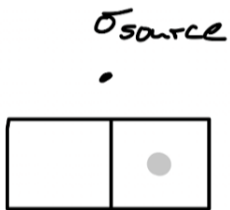
Bath: Simple, Boundless, Translation invariant, Passive

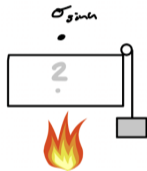
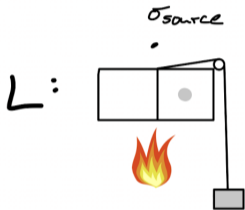


Bath: Simple, Boundless, Translation invariant, Passive

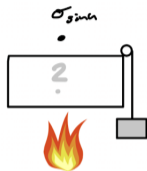
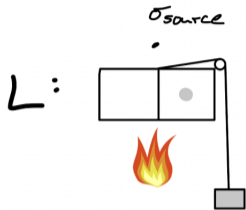


What about just information?



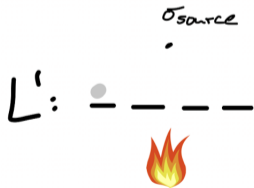


$$W_L^E(p_\ell) = -k_E T \log(2)$$

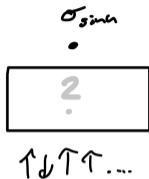
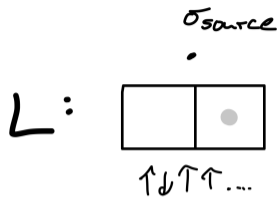


$$W_L^E(p_\ell) = -k_E T \log(2)$$

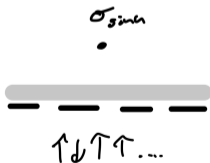
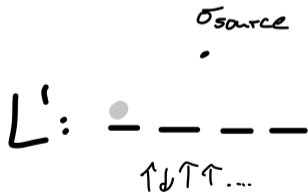
$$\frac{S(L)}{S(L')} = \frac{W_L^E(p_\ell)}{W_{L'}^E(p_{\ell'})}$$



$$W_{L'}^E(p_{\ell'}) = -k_E T \log(4)$$



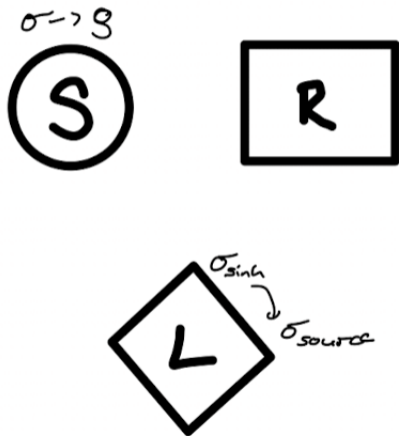
$$W_L^{\text{ang.}}(p_\ell) = -k_{\text{ang.}} T \log(2)$$



$$\frac{S(L)}{S(L')} = \frac{W_L^{\text{ang.}}(p_\ell)}{W_{L'}^{\text{ang.}}(p_{\ell'})}$$

$$W_{L'}^{\text{ang.}}(p_{\ell'}) = -k_{\text{ang.}} T \log(4)$$

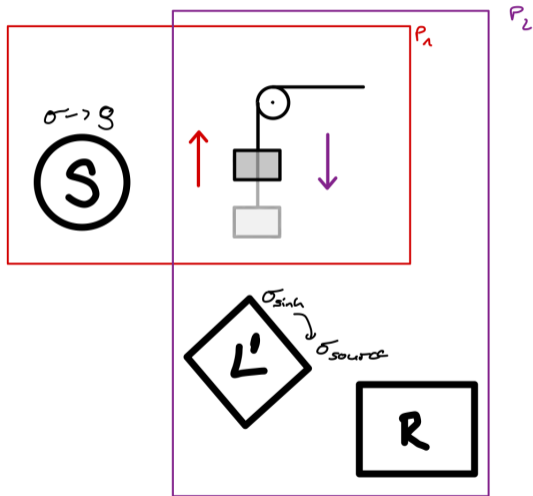
Entropy from Landauer boxes



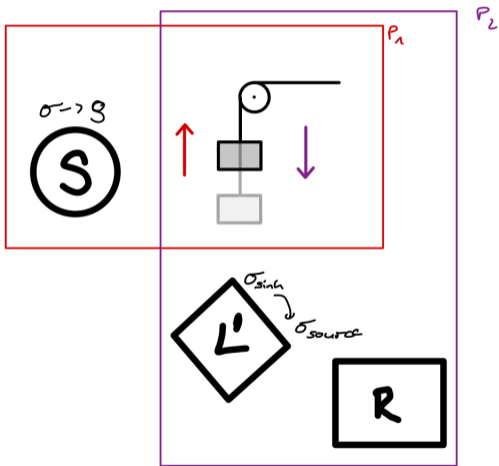
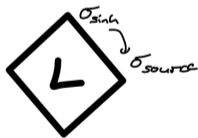
Entropy from Landauer boxes



Entropy from Landauer boxes



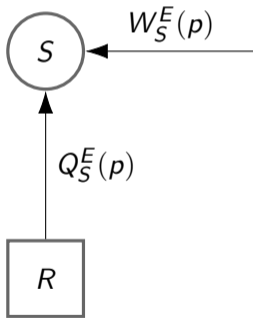
Entropy from Landauer boxes



$$\Delta S(\sigma \rightarrow \rho) = S(L) - S(L')$$

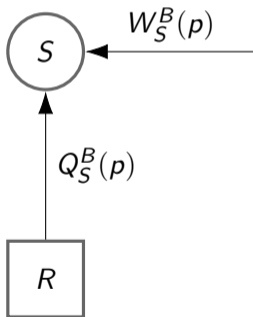
Recovers the usual definition of entropy

$$\Delta S(\rho \rightarrow \sigma) = \frac{Q_S^E(p)}{T_R} \quad \sim \quad dS = \frac{dU - dW}{T}$$



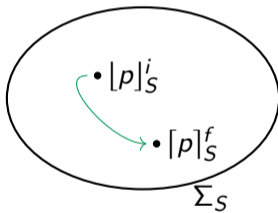
Recovers the usual definition of entropy

$$\Delta S(\rho \rightarrow \sigma) = \sum_B \frac{Q_S^B(\rho)}{k_B T_R^B}$$



Entropy always increases

$$S([p]_S^i \rightarrow [p]_S^f) \geq 0$$



Recover results of thermodynamics

- ▶ Generalized second law $\sum_B \frac{W_S^B(p)}{k_B T_R^B} \geq 0$
- ▶ Generalized Clausius theorem
- ▶ Generalized Carnot theorem
- ▶ Generalized zeroth law

Application: Black holes

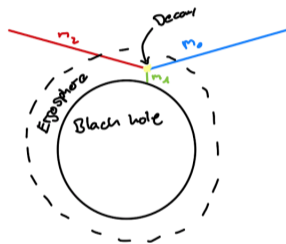
Systems and processes: black hole

- ▶ State space: $\Sigma_{BH} = \{(M, L) \mid M^2 \geq L^2/M^2\}$
- ▶ Dump mass m from infinity



- ▶ Initial state: (M_1, L_1)
- ▶ Final state: $(M_1 + m, L_1)$

- ▶ Reversible Penrose process:



- ▶ Initial state: (M_1, L_1)
- ▶ Final state: $(M_1 + m_1, L_1 + \ell_1)$
- ▶ Irreducible mass is constant

$$M_{1/2}^2 = M_{irr}^2 + \frac{L_{1/2}^2}{4M_{irr}^2}$$

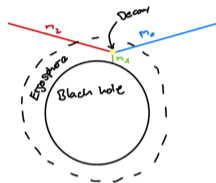
First law: black hole

- ▶ Dump mass m from infinity:



- ▶ Initial state: (M_1, L_1)
- ▶ Final state: $(M_1 + m, L_1)$
- ▶ $W_{E,BH} = m, W_{L,BH} = 0$

- ▶ Reversible Penrose process:

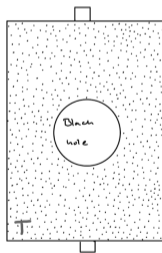


- ▶ Initial state: (M_1, L_1)
- ▶ Final state: $(M_1 + m_1, L_1 + \ell_1)$
- ▶ Irreducible mass is constant
$$M_{1/2}^2 = M_{irr}^2 + \frac{L_{1/2}^2}{4M_{irr}^2}$$
- ▶ $W_{E,BH} = m_1, W_{L,BH} = \ell_1$

$$U_{E,BH} = M + const., U_{L,BH} = L + const.$$

Entropy: black hole

- ▶ We can use the reversible Penrose process
- ▶ Need a process that changes the irreducible mass M_{irr} of the black hole.
- ▶ Use isothermic process considered by Kaburaki and Okamoto²



$$\begin{aligned}\Delta S &= 4\pi (M_{2,irr}^2 - M_{1,irr}^2) \\ &= (A_2 - A_1)/4\end{aligned}$$

²Osamu Kaburaki and Isao Okamoto. "Kerr black holes as a Carnot engine". In: *Physical Review D* 43.2 (1991), pp. 340–345. DOI: 10.1103/physrevd.43.340.

Conclusion

- ▶ Extended the phenomenological thermodynamic framework to multiple quantities.
- ▶ Defined thermodynamic entropy by converting between quantities using Landauer's principle.
- ▶ A black hole fits into the the thermodynamic framework.
- ▶ We find the expected internal energy, angular momentum and entropy using the thermodynamic framework.